

1 Strong vertical mixing of deep water of a stratified reservoir 2 during the Maule earthquake, central Chile (M_w 8.8)

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6 [1] Vertical profiles of water temperature in Rapel reservoir
7 (central Chile, 34° 2' 23" S, 71° 35' 23" W, 110 m.a.s.l.),
8 recorded the hydrodynamic response of the stratified water
9 column of the reservoir produced by the 8.8 magnitude Maule
10 earthquake (epicentre located in 36°15'36"S, 73°14' 20"W),
11 showing an intense vertical mixing of deep water. The ver-
12 tical mixing was characterized by changes in the potential
13 energy of the water column and the eddy diffusivity of the
14 mixing, and these values were used to estimate the mixing
15 efficiency of the event. It is hypothesized that the turbulent
16 flow that mixed the water column was mainly induced by the
17 oscillation of the bed, although surface gravitational waves
18 could also have contributed to the vertical mixing. So far, this
19 is the first time that the hydrodynamic response of a reservoir
20 during a mega-earthquake is measured and described.
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25 1. Introduction

26 [2] On 27 February 2010, a 8.8 magnitude earthquake
27 followed by a tsunami hit Central Chile killing in together
28 more than 500 peoples, and causing damages in the national
29 infrastructure amounted in several thousand of million
30 American dollars [*Madariaga et al.*, 2010; *Kaiser and*
31 *Regalado*, 2010; *Regalado*, 2010; *Lay et al.*, 2010]. Besides
32 the consequences that a mega-earthquake has on the human
33 activity, the released energy during the Maule earthquake and
34 the way on which that release occurs, triggered natural phe-
35 nomena that are difficult to be measured under normal con-
36 ditions [*Montgomery and Manga*, 2003; *Farias et al.*, 2010].
37 One of these events is documented and discussed here, and
38 corresponds to the intense vertical mixing of the stratified
39 Rapel reservoir during the Maule earthquake.

40 2. Methods and Results

41 2.1. Field Measurements and Mixing Characterization

42 [3] Rapel reservoir is an artificial lake located at about
43 120 km from Santiago, the capital city of Chile (see
44 Figure 1a). It was built in 1968 for hydro-power generation

with a maximum water depth of 85 m near the dam 52
[*Contreras et al.*, 1994; *Vila et al.*, 1997; *de la Fuente and* 53
Niño, 2008]; however, reservoir sedimentation raised the 54
bottom elevation in 30 m, being today the maximum water 55
depth of 55 m. Water withdrawals for hydro-power genera- 56
tion are located at about 45 m depth, elevation at which a 57
density interface develops during summer time [*de la Fuente* 58
and Niño, 2008]. 59

[4] During the Maule earthquake occurred on day 58 of 60
2010, a thermistor chain located at about 350 m from the dam 61
(Figure 1b), recorded time-series of water temperature every 62
30 minutes at 11 different depths, with thermistors spaced 63
each other by 5 m, such as the lowest thermistor is located 64
below the density interface. Water depth in point of the 65
thermistor chain location was about 54 m during Maule 66
earthquake. Measurements during the days before and after 67
the earthquake are shown in the contour plot of Figure 2a, in 68
which the intense mixing in the deepest region of the lake 69
(below 25 m depth) is observed during the Maule earthquake 70
that is marked by the vertical dashed line. This vertical mixing 71
is clearer seeing in Figure 2b in which time series of water 72
temperature at 45 and 50 m depth show a strong temperature 73
gradient of about 2°C in 5 m before the earthquake, followed 74
by the nearly complete vertical-homogenization of the water 75
temperature after the earthquake marked by the vertical 76
dashed line. This vertical mixing is also shown in Figure 2c, 77
where both vertical profiles are separated each other by 78
30 minutes, thus showing that vertical mixing was quite 79
instantaneous and localized in deepest areas of the water 80
column, while surface water densities (above 25 m depth) did 81
not change product of the earthquake. 82

[5] In terms of energy budget, changes in water tempera- 83
ture indicate changes in the potential energy of the lake, 84
which increases in case of a mixing event [*Fischer et al.*, 85
1979; *Imberger and Hamblin*, 1982; *Wüest et al.*, 2000]. 86
Potential energy per unit of area of bed is estimated as the 87
gravity centre elevation times the total mass of the water 88
column [*Wüest et al.*, 2000; *Antenucci et al.*, 2000]; therefore, 89
since the deepest waters are denser in stratified conditions, the 90
gravity centre is deeper than in homogenous conditions, thus 91
being required an external source of energy to induce vertical 92
mixing [*Fischer et al.*, 1979; *Imberger and Hamblin*, 1982; 93
Wüest et al., 2000]. In case of Rapel reservoir, the potential 94
energy per unit of bed area of the water column between the 95
elevation of the deepest thermistor (z_t) and the free surface 96
at the elevation H , was computed by solving the integral 97

$$PE = \int_{z_t}^H \rho g z \, dz, \quad (1)$$

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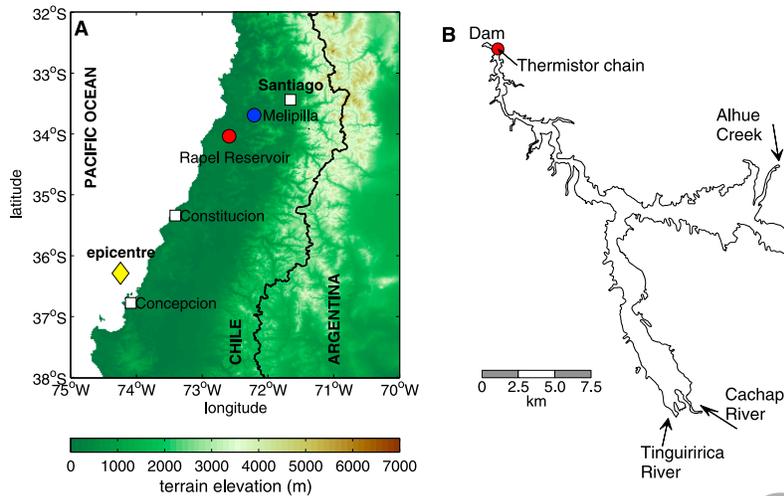


Figure 1. (a) Geographic location of Rapel reservoir, Melipilla seismological station, the earthquake epicentre and mayor cities of the area. (b) Perimeter of Rapel reservoir and arrows indicate inflows (Cachapoal and Tinguiririca rivers and Alhue creek). Red dot specifies the location of thermistor chain.

98 based on a trapezoidal integration, obtaining the time series of
 99 PE shown in Figure 3a, where it is observed that mixing due
 100 to the Maule earthquake increased the potential energy per
 101 area of bed in 0.8 kJm^{-2} . To put in context this change in
 102 potential energy per unit of area of bed, about 1% of the total
 103 energy in the internal wave field of a stratified lake may be
 104 used in raising the potential energy [Gloor *et al.*, 2000;
 105 Shimizu and Imberger, 2008], and the total energy per unit of
 106 area in the internal wave field induced by typhoons passing
 107 over Lake Biwa (Japan) was estimated in 0.52 kJm^{-2}
 108 [Shimizu *et al.*, 2007]. Therefore, typhoons on Lake Biwa
 109 may induce vertical mixing equivalent to 0.005 kJm^{-2} (1% of
 110 the total energy in the flow) in a timescale of days, while
 111 Maule earthquake produced mixing two order of magnitude
 112 larger in a timescale of minutes.
 113 [6] Furthermore, the eddy diffusivity, K_ρ , which char-
 114 acterizes the mixing induced by the Maule earthquake, was

estimated by solving the vertical diffusion equation of the
 water density, ρ ,

$$\frac{\partial \rho}{\partial t} = \frac{\partial}{\partial z} \left(K_\rho \frac{\partial \rho}{\partial z} \right). \quad (2)$$

This equation was numerically solved following Patankar
 [1980], with the initial condition defined by the vertical
 profile measured before the earthquake (red line in Figure 2c),
 and a constant and homogeneous K_ρ was calibrated such as
 obtaining the best fit with measurements 30 minutes after
 (blue line in Figure 2c). Boundary conditions to solve (2)
 were no-flux boundary condition at the bottom of Rapel
 reservoir ($z = 0 \text{ m}$, or 54 m depth during the earthquake), and,
 because of the water density above 25 m depth did not change
 during the earthquake (Figure 2c), $\rho = 997.72 \text{ kg m}^{-3}$ was
 imposed at $z = 29 \text{ m}$ (or 25 m depth). Furthermore, the initial

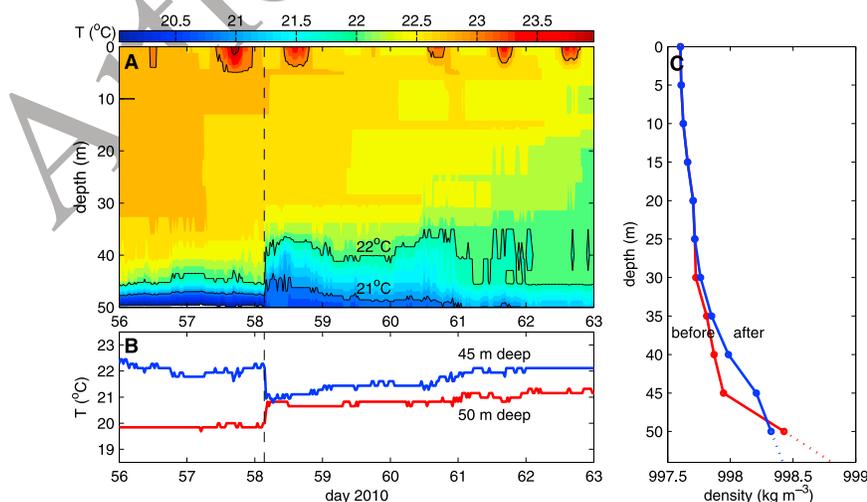


Figure 2. (a) Contour of temporal evolution of vertical profiles of water temperature recorded by the thermistor chain during Maule earthquake (day 58 of 2010). (b) Time series of water temperature measured at 45 m (blue line) and 50 m (red line) depth. Vertical dashed line in Figures 2a and 2b mark the Maule earthquake. (c) Vertical profiles of water density before and after Maule earthquake. Profiles were measured separated each other by 30 min. Dashed line extrapolate measurements toward the bottom of the reservoir at 54 m.

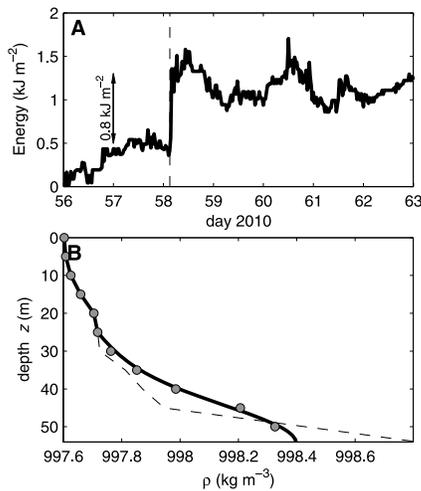


Figure 3. (a) Time series of computed potential energy per unit of area based on water temperature measurements. Vertical dashed line marks the Maule earthquake. (b) Vertical profiles of water density before the earthquake (dashed line), after the earthquake (dots), and simulated after the earthquake for $K_\rho = 2.0 \times 10^{-2} \text{ m}^2 \text{ s}^{-1}$ (solid line).

128 water density below the deepest thermistor was linearly
 129 extrapolated toward the bottom of the lake (see dashed line
 130 in Figure 3b). K_ρ was then calibrated by minimizing the
 131 least square error between measured and simulated water
 132 density, obtaining $K_\rho = 2.0 \times 10^{-2} \text{ m}^2 \text{ s}^{-1}$ and the corre-
 133 sponding profile is shown in Figure 3b. Values for K_ρ in
 134 normal conditions without an earthquake varies between 10^{-4}
 135 and $10^{-3} \text{ m}^2 \text{ s}^{-1}$ [MacIntyre et al., 1999; Lorke et al., 2002;
 136 Wüest and Lorke, 2003].

137 2.2. Mixing Efficiency

138 [7] Two different approaches are used to estimate the
 139 energy per unit of bed area that was delivered by the earth-
 140 quake to Rapel reservoir: the second problem of Stokes
 141 [Batchelor, 1967; Lorke et al., 2002], and Shih et al. [2005]
 142 and Ivey et al. [2008] parameterization of mixing in strati-
 143 fied flows.

144 [8] The second problem of Stokes consist in solving the
 145 horizontal velocity parallel to a wall, $u(z,t)$, with z the vertical
 146 coordinate and t the time, which is the result of inducing
 147 the flow by an oscillatory boundary conditions at the base of
 148 $u(0,t) = u_o \cos(\omega t)$, where u_o is the amplitude of the oscillation
 149 and ω the frequency at which occurs [Batchelor, 1967; Lorke
 150 et al., 2002]. In this way, shear stresses induces an oscillatory
 151 flow in a region characterized by the thickness of the Stokes
 152 layer, $d = 2\pi\sqrt{2\nu/\omega}$, above which no horizontal velocities
 153 are expected, with ν denoting the water viscosity that is
 154 considered to be the turbulent eddy viscosity in the benthic
 155 boundary layer [Lorke et al., 2002], and is assumed equal to
 156 $K_\rho = 2.0 \times 10^{-2} \text{ m}^2 \text{ s}^{-1}$ since the turbulent Prandtl number is
 157 close to 1.0 [Rodi, 1983; Pope, 2000]. Using the velocity
 158 profile of the Stokes problem, the work per unit of area of bed,
 159 per oscillating period done by the oscillating on the flow is

$$w = \int_0^{2\pi/\omega} \rho\nu \frac{\partial u}{\partial z} \Big|_{z=0} u_o \cos(\omega t) dt = \rho u_o^2 \pi \sqrt{\frac{\nu}{2\omega}}. \quad (3)$$

Then, the net energy per unit of area of bed for a mono- 160
 chromatic oscillation, is calculated as the work per period 161
 (w , equation (3)) times the number of periods contained on 162
 a time T_s , giving: 163

$$W = \rho u_o^2 T_s \sqrt{\frac{\nu\omega}{8}}. \quad (4)$$

In this way, the net work per unit of area of bed done by the 164
 earthquake on the reservoir, W , can be contrasted with the 165
 changes in the potential energy of Rapel reservoir produced 166
 by the Maule earthquake ($\Delta PE = 0.8 \text{ kJ m}^{-2}$, Figure 3a), for 167
 which it is required knowing T_s , ν , u_o and ω . T_s is assumed 168
 equal to the duration of the earthquake (90 s), and $\nu = K_\rho$ as it 169
 was previously discussed. The other two variables, u_o and ω , 170
 are taken from the accelerometer station located in Melipilla 171
 station (see Figure 1a), which is the closest available seis- 172
 mological station of Rapel reservoir (Seismological Service 173
 of Universidad de Chile, 2010, <http://ssn.dgf.uchile.cl>). Land 174
 velocities along west–east and south–north axes were com- 175
 puted by integrating in time the accelerations recorded in 176
 Melipilla, and long waves were filtered out [Boore and 177
 Bommer, 2005] by subtracting to the original time series 178
 the moving average of the velocity in a time window of 1s. As 179
 an example, the resulting time series of ground velocity along 180
 the east–west axis is shown in Figure 4a. Furthermore, the 181
 energy spectrum of each component of the land velocity was 182
 computed, and the specific kinetic energy of the earthquake 183
 per unit of frequency, $u_o^2(\omega)$ used in (4), was computed by 184
 adding the energy spectrums of the east–west and south– 185
 north components of the horizontal land velocity. Then, the 186
 exchanged work per unit of frequency per unit of bed area, 187
 $W(\omega)$, was computed based on (4), and the corresponding 188
 spectral form is shown in Figure 4b ($\nu = K_\rho = 2.0 \times 10^{-2} \text{ m}^2$ 189
 s^{-1} and $T_s = 90 \text{ s}$). Finally, the total energy per unit of area of 190
 bed that entered to the reservoir (ΔE) was computed by 191
 integrating $W(\omega)$ (Figure 4b) in the range of frequencies of the 192
 earthquake (between 10^{-1} and 10^2 rad s^{-1} , Figure 4b), 193
 obtaining $\Delta E = 25.6 \text{ kJ m}^{-2}$ that is about 32 times larger than 194
 ΔPE , which is equivalent to a mixing efficiency of $\gamma_{mix} =$ 195
 $\Delta PE(\Delta E)^{-1} = 3.2\%$. It worth noticing that if this estimation 196

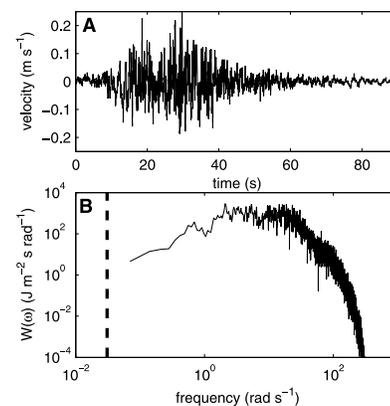


Figure 4. (a) Computed time series of land velocity in Melipilla seismological station, along the West–East direc-
 tion. (b) Estimated spectral energy per unit of frequency
 delivered by the earthquake to the reservoir. Vertical dashed
 line defines maximum buoyancy frequency before the
 earthquake.

197 is done by also considering vertical land velocities $\Delta E =$
198 28.3 kJ m^{-2} , thus showing that in the first approximation the
199 vertical accelerations can be neglected.

200 [9] The second way of computing the mixing efficiency,
201 γ_{mix} , is based on the results of *Shih et al.* [2005] and *Ivey et al.*
202 [2008]. These authors showed that the energetic regime is
203 described by

$$\frac{K_\rho}{\nu_m} = 2 \left(\frac{\varepsilon}{\nu_m N^2} \right)^{0.5} = \gamma_{mix} \left(\frac{\varepsilon}{\nu_m N^2} \right), \quad (5)$$

204 where ν_m is the molecular kinematic viscosity ($1.3 \times 10^{-6} \text{ m}^2$
205 s^{-1}), N is the buoyancy frequency (rad s^{-1}), and ε is the rate of
206 TKE dissipation ($\text{m}^2 \text{ s}^{-3}$). Then, using $K_\rho = 2.0 \times 10^{-2} \text{ m}^2 \text{ s}^{-1}$,

$$\gamma_{mix} = 4 \left(\frac{K_\rho}{\nu_m} \right)^{-1} = 0.03\%. \quad (6)$$

207 and $\Delta E = \Delta PE \gamma_{mix}^{-1} = 3000 \text{ kJm}^{-2}$ which is equivalent to 24%
208 for a 50 m height water column.

209 3. Discussion

210 [10] In this article, we analyzed unique measurements
211 of the hydrodynamic response of Rapel reservoir due to the
212 8.8 M_w earthquake that hit Central Chile in February 2010.
213 Based on these measurements, the mixing was character-
214 ized by a change in the potential energy per unit of area of
215 bed, $\Delta PE = 0.8 \text{ kJm}^{-2}$, with an eddy diffusivity of $K_\rho = 2.0 \times$
216 $10^{-2} \text{ m}^2 \text{ s}^{-1}$, and a mixing efficiency that varied between 0.03
217 to 3%. Although changes in the potential energy per unit of
218 area of the bed were calculated from the measurements of the
219 thermistor chain, several assumptions were required to esti-
220 mate both the eddy diffusivity and mixing efficiency.

221 [11] First of all, to estimate the eddy diffusivity it was
222 required to define the timescale in which mixing occurred,
223 which is something in between the earthquake duration (90 s)
224 and 30 min that is the time-step of the measurements. Pre-
225 vious studies showed that diffusion timescale of turbulent
226 patches $4N^{-1}$ [*Fernando*, 1988] is equal to 13.6 min, when a
227 vertically averaged N is used. However, this timescale is valid
228 for three-dimensional turbulent patches, while the turbulence
229 induced by the earthquake is bounded by the bottom of the
230 reservoir.

231 [12] Second, it was assumed that the vertical mixing was
232 driven by the turbulent flow induced in the benthic boundary
233 layer by the oscillation of the bed. Questions arise whether or
234 not this was the only mixing mechanism excited by the
235 earthquake. The excitation of surface and internal waves is a
236 plausible hypothesis for also explaining the observed increase
237 in vertical mixing, as the benthic boundary is usually ener-
238 gized by shear-induced turbulence or by internal wave
239 breaking [*Gloor et al.*, 2000; *Boegman et al.*, 2003; *Lorke*
240 *et al.*, 2005]. To preliminarily define whether the excitation
241 of surface and internal waves by the earthquake is possible, it
242 is considered that only those waves with natural frequencies
243 near the earthquake frequency are likely to be excited, thus
244 neglecting nonlinear energy transfers among waves as a
245 mechanism of excitation [*de la Fuente et al.*, 2010]. With this
246 consideration in mind, internal waves are not possible to be
247 excited because the oscillations induced by the earthquake
248 delivered energy in flow-scales of $\omega = 10 \text{ rad s}^{-1}$, which are
249 faster than the buoyancy frequency (vertical line in Figure 4b)

that defines the higher frequency at which a stratified fluid
oscillates [*Boegman et al.*, 2003]. Surface waves, in contrast,
are much faster than the internal waves, thus being more
likely to be excited by the land oscillations of the earthquake.
In this case, the wavelength, L , of a gravitational surface
waves excited by an oscillating forcing can be estimated as

$$L = \frac{2\pi}{\omega} c \quad (7)$$

where $c = \sqrt{gh}$ is the surface wave celerity of a water column
with a depth h . With $h = 54 \text{ m}$ and $\omega = 10 \text{ rad s}^{-1}$, the
wavelength is $L = 14.5 \text{ m}$, which are small waves rather
than basin-scale waves. However, because available mea-
surements are every 30 minutes, it is unclear if surface
waves were excited by the earthquake, and how much energy
flood into them.

[13] Third, the mixing efficiency, γ_{mix} , was estimated with
the second problem of Stokes, obtaining $\gamma_{mix} = 3.2\%$; and
with *Shih et al.* [2005] and *Ivey et al.* [2008] parameterization
of mixing for stratified flows, obtaining $\gamma_{mix} = 0.03\%$, which
have a difference of two-order of magnitude. The problem is
that neither of these two approaches is actually valid for the
conditions of the earthquake. On the one hand, the Stokes
solution is valid for a constant and homogeneous ν . On the
other hand, equations (5) and (6) are valid for $\varepsilon(\nu_m N^2)^{-2} <$
 10^6 , whereas during the earthquake $\varepsilon(\nu_m N^2)^{-1} \sim O(10^8)$.
Further research is needed to elucidate this issue.

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